

Conservation Law Based Updating Schemes for Overset Meshes

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Conservation Law Based Update Schemes

- In this talk a method for the exchange of data at overset mesh interfaces is outlined
- The method is based on enforcing the conservation laws on a cell-wise basis rather than the traditional approaches involving interpolation
- Reexamine some old questions in a new light
- Consider the following conservation law

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} = 0$$

Conservation Law Based Update Schemes

- In order to solve this equation numerically on a single grid, the domain is partitioned into discrete cells $j = \left[x_j - \frac{\Delta x}{2}, x_j + \frac{\Delta x}{2} \right]$

- The variable q is averaged over cell j as follows

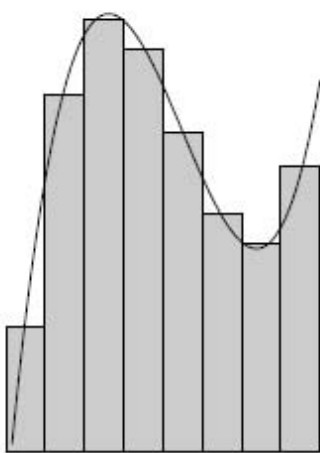
$$q_j^n = \int_{x_j - \frac{\Delta x}{2}}^{x_j + \frac{\Delta x}{2}} q(x, t_n) dx ; t_n = n \Delta t ; \text{etc.}$$

- The value of the cell averaged variable q_j in cell j is updated as follows

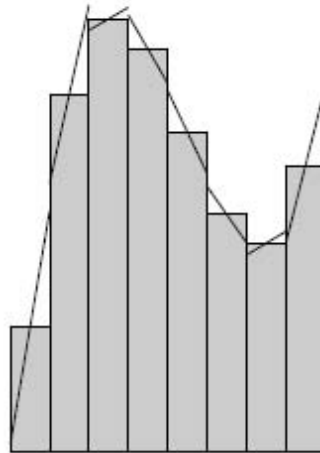
$$q_j^{n+1} = q_j^n - \frac{\Delta t}{\Delta x} \left(F(q_l_{j+1}, q_r_{j+1}) - F(q_l_j, q_r_j) \right)$$

Conservation Law Based Update Schemes

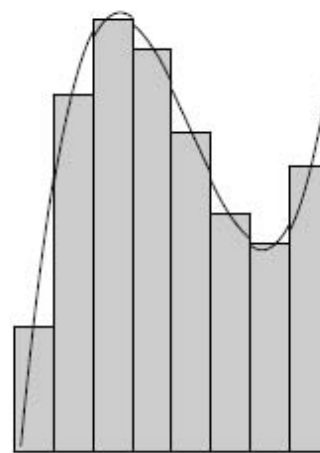
- Here q_l, q_r are the values of q reconstructed from the cell averages q_j^n and evaluated at the cell interfaces using the MUSCL procedure and F is a numerical flux function (Roe, HLLE++, etc.)



a. Cell averaging of quartic data



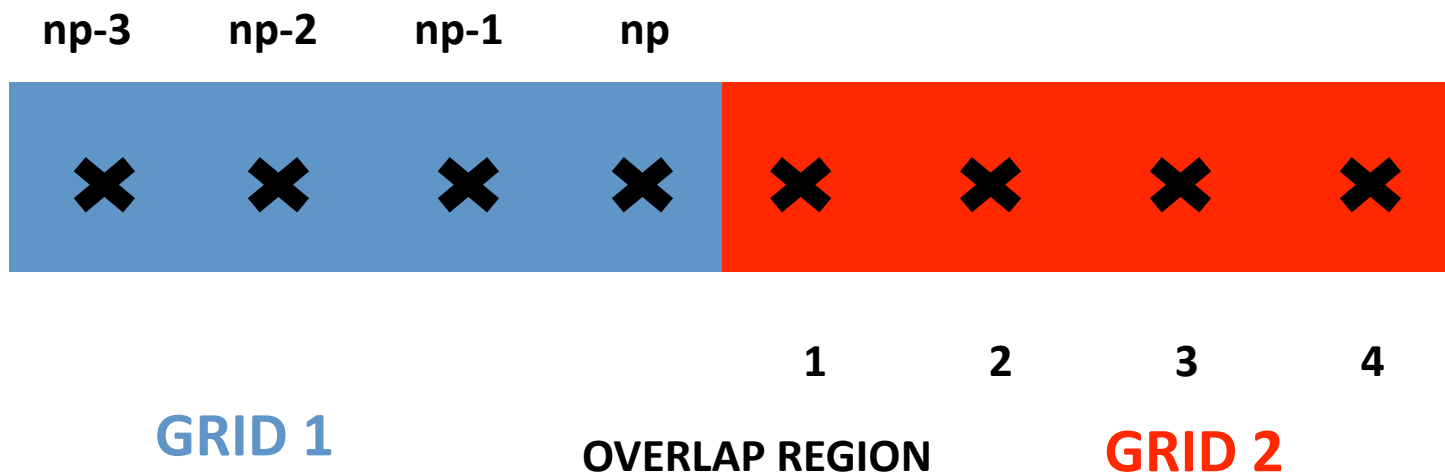
b. Linear reconstruction



c. Quadratic reconstruction (Barth and Ohlberger, 2004)

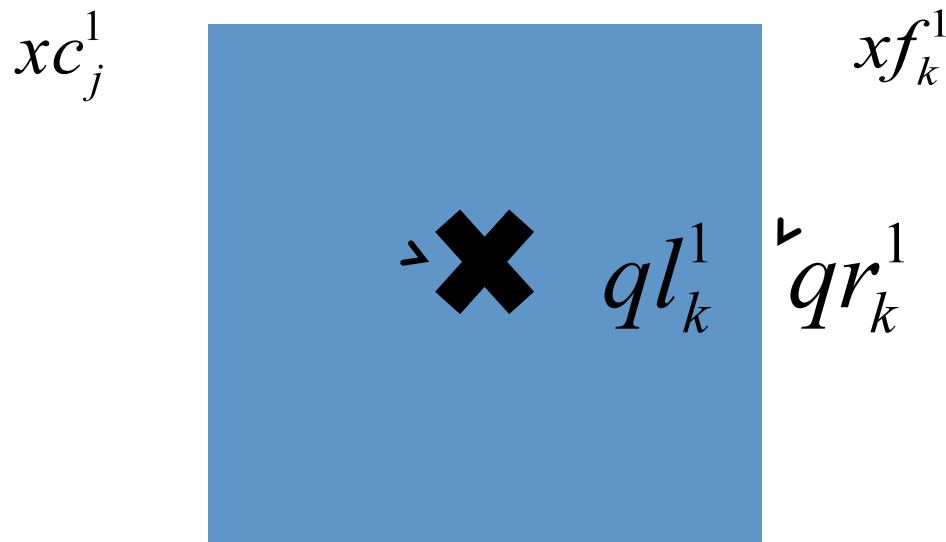
Conservation Law Based Update Schemes

- What is the “best” way to extend this approach to cases of multiple overlapping grids?
- Rather than interpolate variables or fluxes seek alternate formulation
- Use data from both zones to reconstruct interface states
- Consider two grids with some small amount of overlap



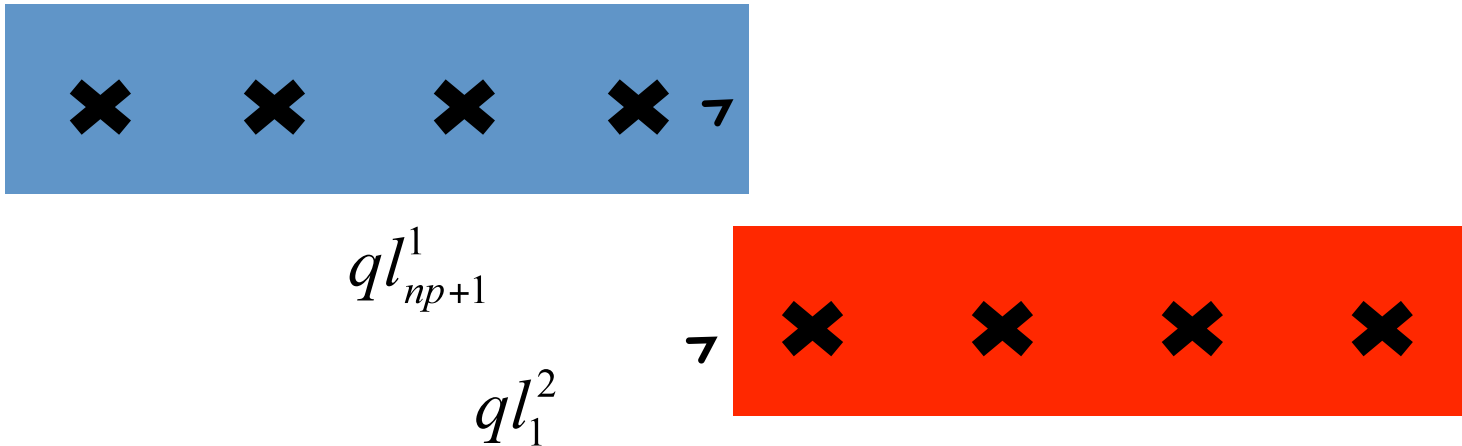
Conservation Law Based Update Schemes

- Let the value of the cell averaged variable in cell j of grid 1 be denoted as q_j^1 , and the cell centroid value of cell j be denoted as xc_j^1 etc
- Let the left and right interface states at face k of grid 1 be denoted as ql_k^1 and qr_k^1 respectively
- Let the face centroid of face k of grid 1 be denoted as xf_k^1 respectively



Conservation Law Based Update Schemes

- Consider computing the cell interface states ql_{np+1}^1 and ql_1^2 as follows



- The states are computed by constructing a least squares monotonically limited approximation to the solution gradient in cell np from mesh 1 using the cells below and then extrapolating q to the cell faces using this gradient

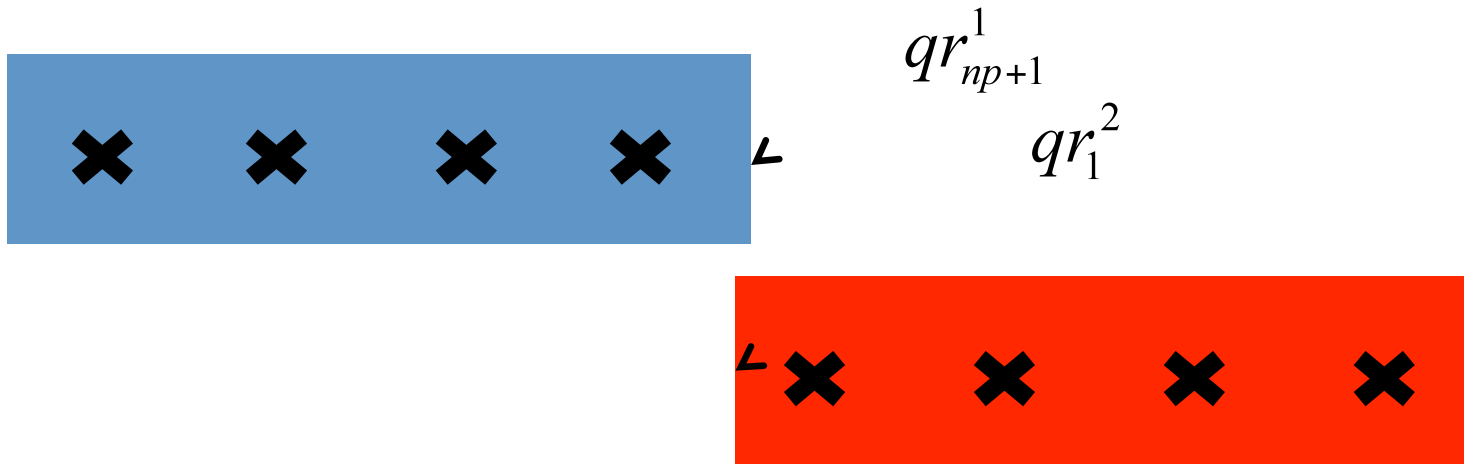
The diagram shows three cells labeled $np-1$, np , and 1 . The first two cells are blue and the third is red. Each cell contains a black 'x' mark.

$$ql_{np+1}^1 = q_{np}^1 + \overline{\nabla} q_{np}^1 \cdot (xf_{np+1}^1 - xc_{np}^1)$$

$$ql_1^2 = q_{np}^1 + \overline{\nabla} q_{np}^1 \cdot (xf_1^2 - xc_{np}^1)$$

Conservation Law Based Update Schemes

- Consider computing the cell interface states qr_{np+1}^1 and qr_1^2 as follows



- The states are computed by constructing a least squares monotonically limited approximation to the solution gradient in cell 1 from mesh 2 using the cells below and then extrapolating q to the cell faces using this gradient

np	1	2	
x	x	x	$qr_{np+1}^1 = q_1^2 + \overline{\nabla} q_1^2 \cdot (xf_{np+1}^1 - xc_1^2)$ $qr_1^2 = q_1^2 + \overline{\nabla} q_1^2 \cdot (xf_1^2 - xc_1^2)$

Conservation Law Based Update Schemes

- A consistent and convergent scheme should be achieved as grid refinement is performed if the distance between the cells centers used in the reconstructions approach zero under grid refinement, i.e.

$$ql_{np+1}^1 \rightarrow ql_1^2 \text{ and } qr_{np+1}^1 \rightarrow qr_1^2 \text{ as } \Delta x \rightarrow 0$$

in that

$$\left\| qr_{np+1}^1 - qr_1^2 \right\| = \left\| \bar{\nabla} q_1^2 \cdot \left(xf_{np+1}^1 - xf_1^2 \right) \right\| \leq \left\| \bar{\nabla} q_1^2 \right\| \left\| \left(xf_{np+1}^1 - xf_1^2 \right) \right\| \leq TV(q) \Delta x$$

etc.

- Here TV is the total variation
- A similar argument holds for meshes with small gaps between the grids
- The basic argument remains unchanged in the case of general overlap

Generalized Lax-Wendroff Theorem

- Recall that a function $q(x,t)$ is considered to be a weak solution to the conservation law

$$\frac{\partial q}{\partial t} + \frac{\partial f(q)}{\partial x} = 0$$

If for any compact, differentiable function $\phi(x,t)$ the following integrals are satisfied

$$\int_0^{\infty} \int_{-\infty}^{\infty} [\phi_t q + \phi_x f(q)] dx dt = - \int_{-\infty}^{\infty} \phi(x,0) q(x,0) dx$$

Generalized Lax-Wendroff Theorem

- In order to demonstrate that these integrals are satisfied for the proposed method, first multiply the cell average update equation by $\phi(x_j^1, t_n)$ for mesh 1, etc.

$$\phi(x_j^1, t_n) (q_j^{1,n+1} - q_j^{1,n}) = -\frac{\Delta t}{\Delta x} \phi(x_j^1, t_n) (F(q_{j+1}^1, q_{j+1}^1) - F(q_j^1, q_j^1))$$

- Next sum over j and n for each mesh as follows

$$\begin{aligned} \Delta t \Delta x \left\{ \sum_{j=-\infty}^{np} \sum_{n=0}^{\infty} \phi(x_j^1, t_n) \frac{(q_j^{1,n+1} - q_j^{1,n})}{\Delta t} + \sum_{j=1}^{\infty} \sum_{n=0}^{\infty} \phi(x_j^2, t_n) \frac{(q_j^{2,n+1} - q_j^{2,n})}{\Delta t} \right. \\ \left. - \sum_{j=-\infty}^{np} \sum_{n=0}^{\infty} \phi(x_j^1, t_n) \frac{(F(q_{j+1}^1, q_{j+1}^1) - F(q_j^1, q_j^1))}{\Delta x} \right. \\ \left. - \sum_{j=1}^{\infty} \sum_{n=0}^{\infty} \phi(x_j^2, t_n) \frac{(F(q_{j+1}^2, q_{j+1}^2) - F(q_j^2, q_j^2))}{\Delta x} \right\} = \end{aligned}$$

Generalized Lax-Wendroff Theorem

- The following summation by parts formulas are useful for the proving convergence of the proposed scheme to a weak solution

$$\sum_{i=0}^N a_j (b_j - b_{j-1}) = a_N b_N - a_0 b_{-1} - \sum_{i=0}^{N-1} b_j (a_{j+1} - a_j) \quad \text{SPB1}$$

$$\sum_{i=0}^N a_j (b_{j+1} - b_j) = a_N b_{N+1} - a_0 b_0 - \sum_{i=1}^N b_j (a_j - a_{j-1}) \quad \text{SPB2}$$

Generalized Lax-Wendroff Theorem

- For the LHS use the SPB2 formula on the n sum

$$\begin{aligned}
 & - \sum_{j=-\infty}^{np} \Delta x \phi(x_j^1, t_0) q_j^{1,0} - \sum_{j=1}^{\infty} \Delta x \phi(x_j^2, t_0) q_j^{2,0} \\
 & - \sum_{n=1}^{\infty} \Delta t \left(\sum_{j=-\infty}^{np} \Delta x q_j^{1,n} \left(\frac{\phi(x_j^1, t_n) - \phi(x_j^1, t_{n-1})}{\Delta t} \right) + \sum_{j=1}^{\infty} \Delta x q_j^{2,n} \left(\frac{\phi(x_j^2, t_n) - \phi(x_j^2, t_{n-1})}{\Delta t} \right) \right)
 \end{aligned}$$

- Note that these terms are discrete approximations to the following integrals

$$- \int_{-\infty}^{xf_{np+1}^1} \phi(x,0) q(x,0) dx - \int_{xf_1^2}^{\infty} \phi(x,0) q(x,0) dx - \int_0^{\infty} \left(\int_{-\infty}^{xf_{np+1}^1} \phi_t q + \int_{xf_1^2}^{\infty} \phi_t q \right) dx dt$$

Generalized Lax-Wendroff Theorem

- For the RHS use the SPB1 formula on the j sum

$$\sum_{n=0}^{\infty} \Delta t \left(\sum_{j=-\infty}^{np} \Delta x F \left(ql_j^1, qr_j^1 \right) \left(\frac{\phi(x_j^1, t_n) - \phi(x_{j-1}^1, t_n)}{\Delta x} \right) + \sum_{j=2}^{\infty} \Delta x F \left(ql_j^2, qr_j^2 \right) \left(\frac{\phi(x_j^2, t_n) - \phi(x_{j-1}^2, t_n)}{\Delta x} \right) \right)$$

$$- \sum_{n=0}^{\infty} \Delta t \Delta x \left(F \left(ql_{np+1}^1, qr_{np+1}^1 \right) \frac{\phi(x_{np}^1, t_n)}{\Delta x} - F \left(ql_1^2, qr_1^2 \right) \frac{\phi(x_1^2, t_n)}{\Delta x} \right)$$

End Terms from SBP Formula

Generalized Lax-Wendroff Theorem

- The end terms at the grid interface become

$$\begin{aligned}
 & - \sum_{n=0}^{\infty} \Delta t \Delta x \left(F \left(ql_{np+1}^1, qr_{np+1}^1 \right) \frac{\phi(x_{np}^1, t_n)}{\Delta x} - F \left(ql_1^2, qr_1^2 \right) \frac{\phi(x_1^2, t_n)}{\Delta x} \right) = \\
 & \sum_{n=0}^{\infty} \Delta t \Delta x \left(F \left(ql_1^2, qr_1^2 \right) \left(\frac{\phi(x_1^2, t_n) - \phi(x_{np}^1, t_n)}{\Delta x} \right) + \frac{\Delta F_{12} \phi(x_{np}^1, t_n)}{\Delta x} \right)
 \end{aligned}$$

where $\Delta F_{12} = F \left(ql_1^2, qr_1^2 \right) - F \left(ql_{np+1}^1, qr_{np+1}^1 \right)$

Generalized Lax-Wendroff Theorem

- We can show that $\Delta F_{12} \rightarrow 0$ as $\Delta x \rightarrow 0$ as follows

$$\Delta F_{12} = F(ql_1^2, qr_1^2) - F(ql_{np+1}^1, qr_{np+1}^1)$$

$$\Delta F_{12} = F(ql_1^2, qr_1^2) - F(ql_1^2 + (ql_{np+1}^1 - ql_1^2), qr_1^2 + (qr_{np+1}^1 - qr_1^2))$$

$$\|\Delta F_{12}\| \leq C \max(\|ql_{np+1}^1 - ql_1^2\|, \|qr_{np+1}^1 - qr_1^2\|) \quad \mathbf{F \text{ is Lipschitz Continuous}}$$

- Given that

$$\|qr_{np+1}^1 - qr_1^2\| = \|\bar{\nabla} q_1^2 \cdot (xf_{np+1}^1 - xf_1^2)\| \leq \|\bar{\nabla} q_1^2\| \| (xf_{np+1}^1 - xf_1^2) \| \leq TV(q)\Delta x, \text{ etc.}$$

it follows that

$\Rightarrow 0$ as $\Delta x \Rightarrow 0$

$$\|\Delta F_{12} \phi(x_{np}^1, t_n)\| \leq C TV(q^n) \Delta x \|\phi(x_{np}^1, t_n)\|$$

Generalized Lax-Wendroff Theorem

- Note that these terms are discrete approximations to the following integrals

$$\int_0^{\infty} \int_{-x f_{np+1}^1}^{\infty} f(q) \phi_x dx dt + \int_0^{\infty} \int_{x f_1^2}^{\infty} f(q) \phi_x dx dt$$

- If $\|x f_{np+1}^1 - x f_1^2\| \rightarrow 0$ as $\Delta x \rightarrow 0$ then terms like

$$- \int_{-\infty}^{x f_{np+1}^1} \phi(x,0) q(x,0) dx - \int_{x f_1^2}^{\infty} \phi(x,0) q(x,0) dx - \int_0^{\infty} \left(\int_{-\infty}^{x f_{np+1}^1} \phi_t q + \int_{x f_1^2}^{\infty} \phi_t q \right) dx dt$$

$$\text{become } - \int_{-\infty}^{\infty} \phi(x,0) q(x,0) dx - \int_0^{\infty} \int_{-\infty}^{\infty} \phi_t q dx dt$$

and the integral conservation law is satisfied under the same caveats invoked for the single grid case

Observations Based on Lax-Wendroff Theorem (1)

- If standard linear interpolation methods are used then no GENERAL bounds may be established a priori for the ΔF_{12} term
- We end up with a source term in the integral conservation law which **MAY** not vanish under grid refinement in the presence of discontinuous solutions.

$$\approx \int_0^{\infty} \Delta F_{12}(t) \varphi(x_R^1, t) dt$$

- These source terms become ODEs along characteristics and may globally pollute the solution
 - Corrupt convergence rates
 - Generate spurious waves

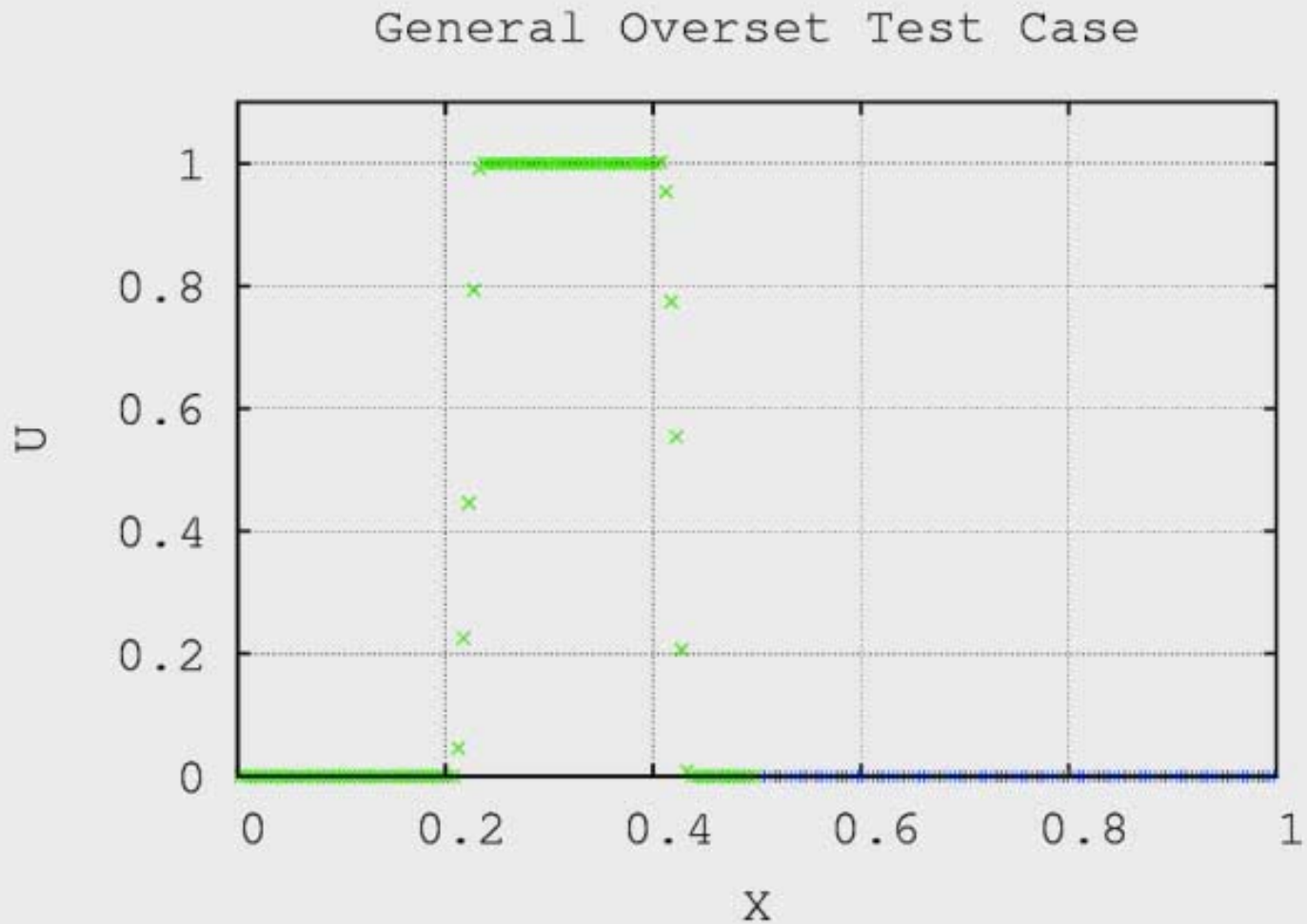
Observations Based on Lax-Wendroff Theorem (2)

- These terms can and will corrupt solutions when
 - Severe mismatches in cell sizes exist
 - Strong discontinuities exist in the grid interface region
 - Slow moving discontinuities exist in the grid interface region
 - ***All three exist simultaneously!***
- These facts need to be taken into consideration when designing grid “surgery” algorithms
 - helps explain the success of the “LEVEL2” interpolation approach used in PEGASUS5 and other grid connectivity codes

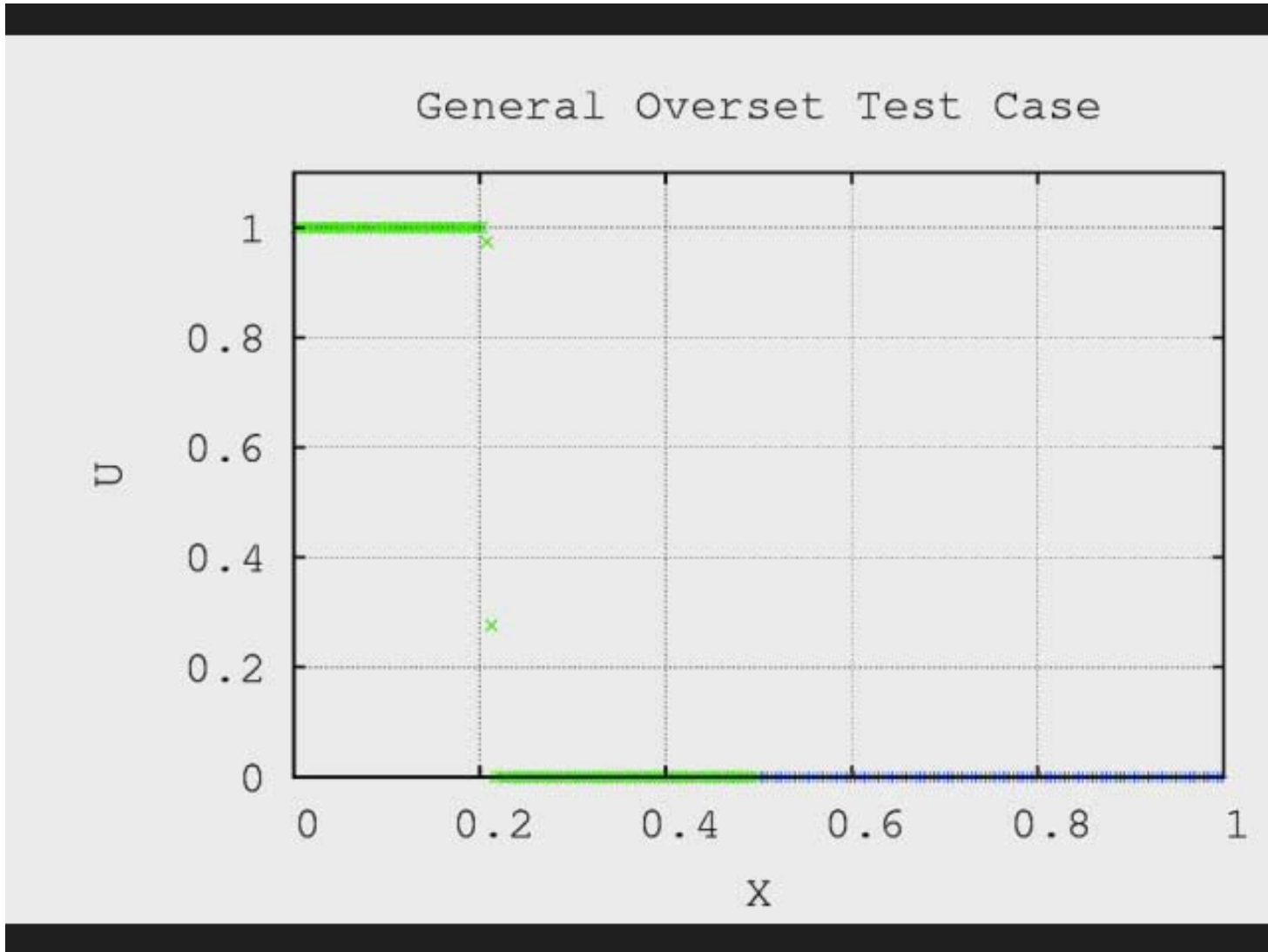
Examples Cases

- SIMPLE ADVECTION (LINEAR DISCONTINUITY)
- BURGERS EQUATION (NONLINEAR DISCONTINUITY)
- EULER EQUATION
 - SLOW MOVING SHOCK (use HLLE++ Riemann Solver)
- All these cases have a small gap in between the grids!
- Use Barth-Jespersen Limiter

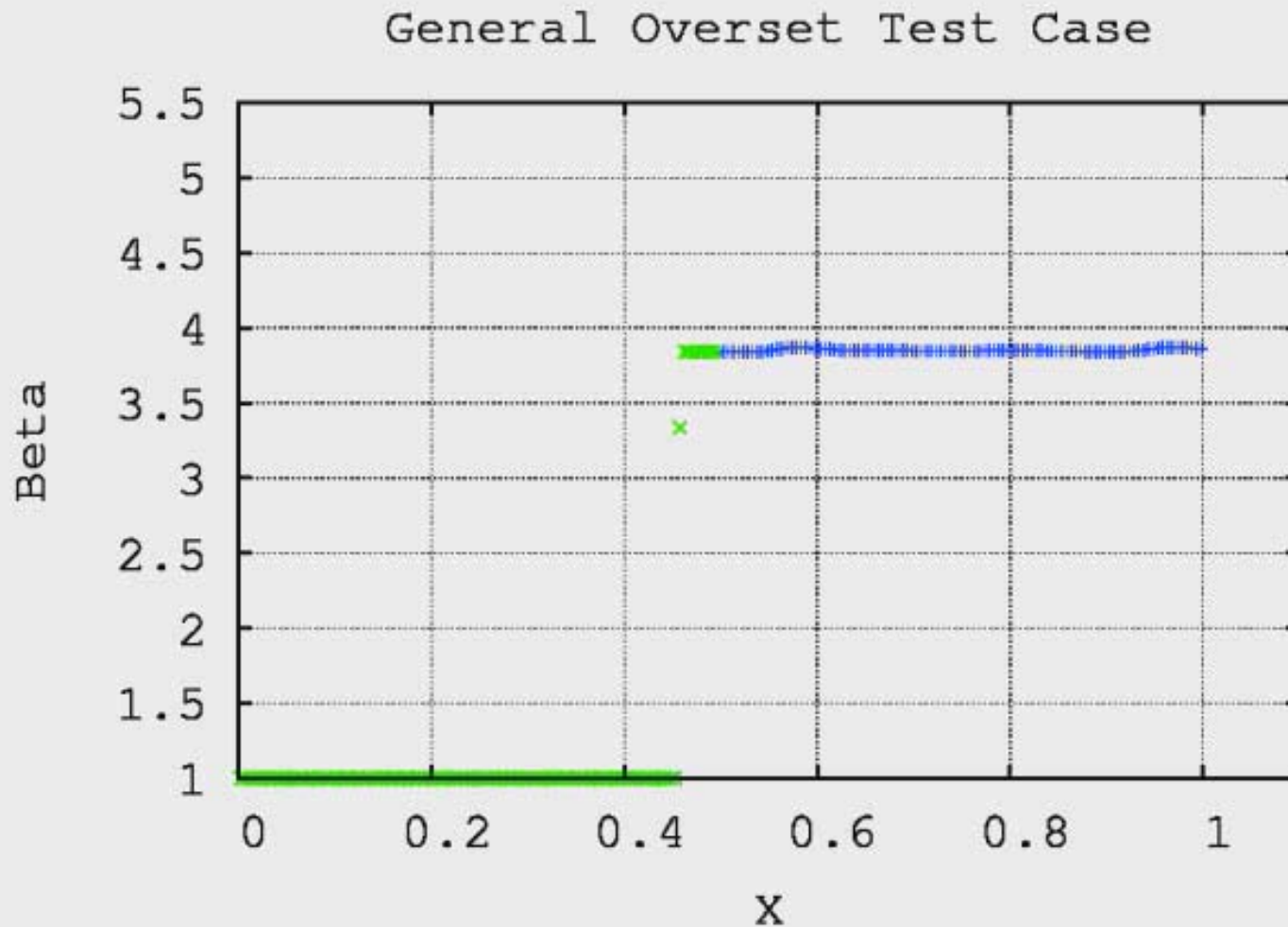
Simple Advection



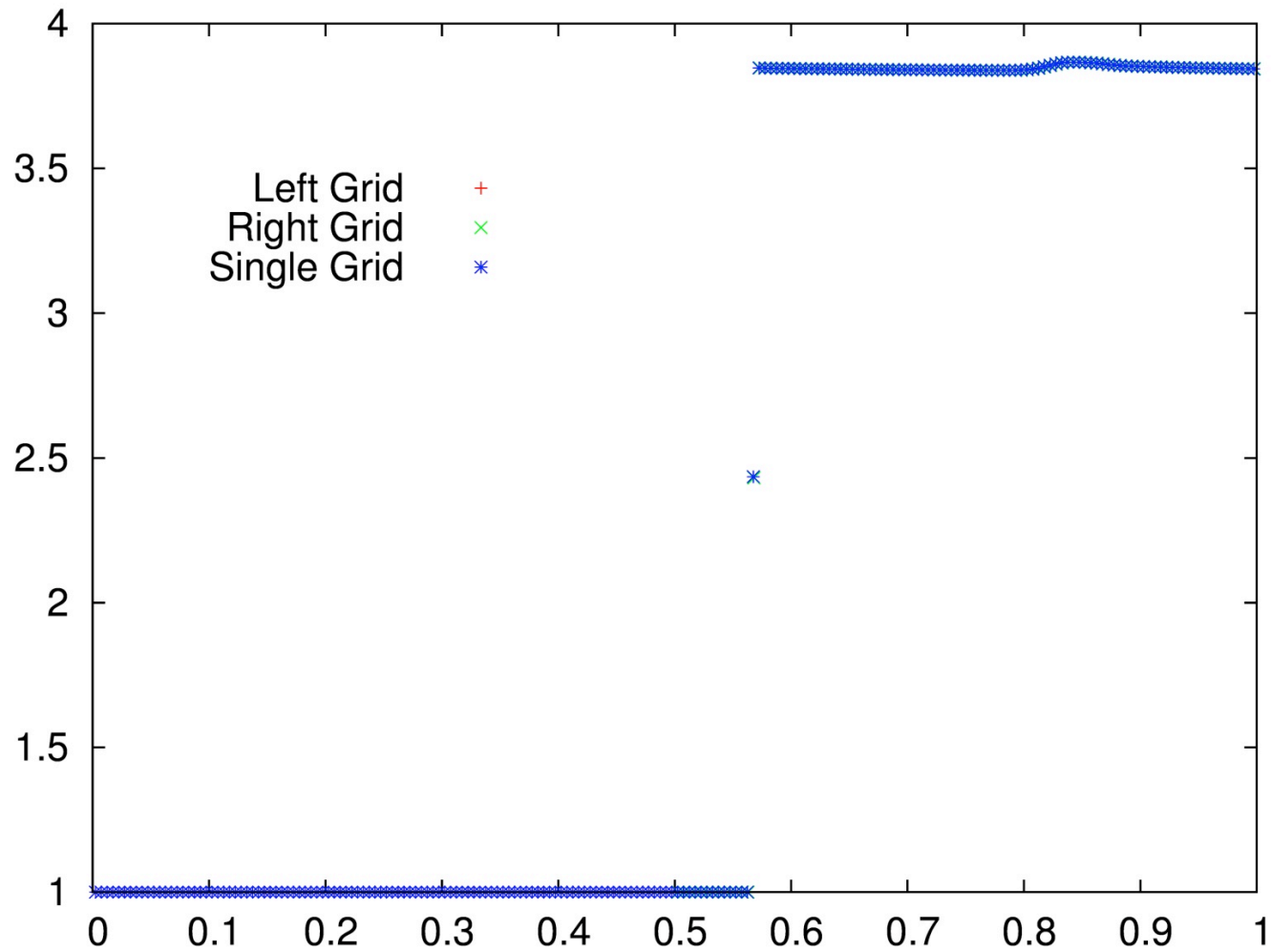
Burgers Equation



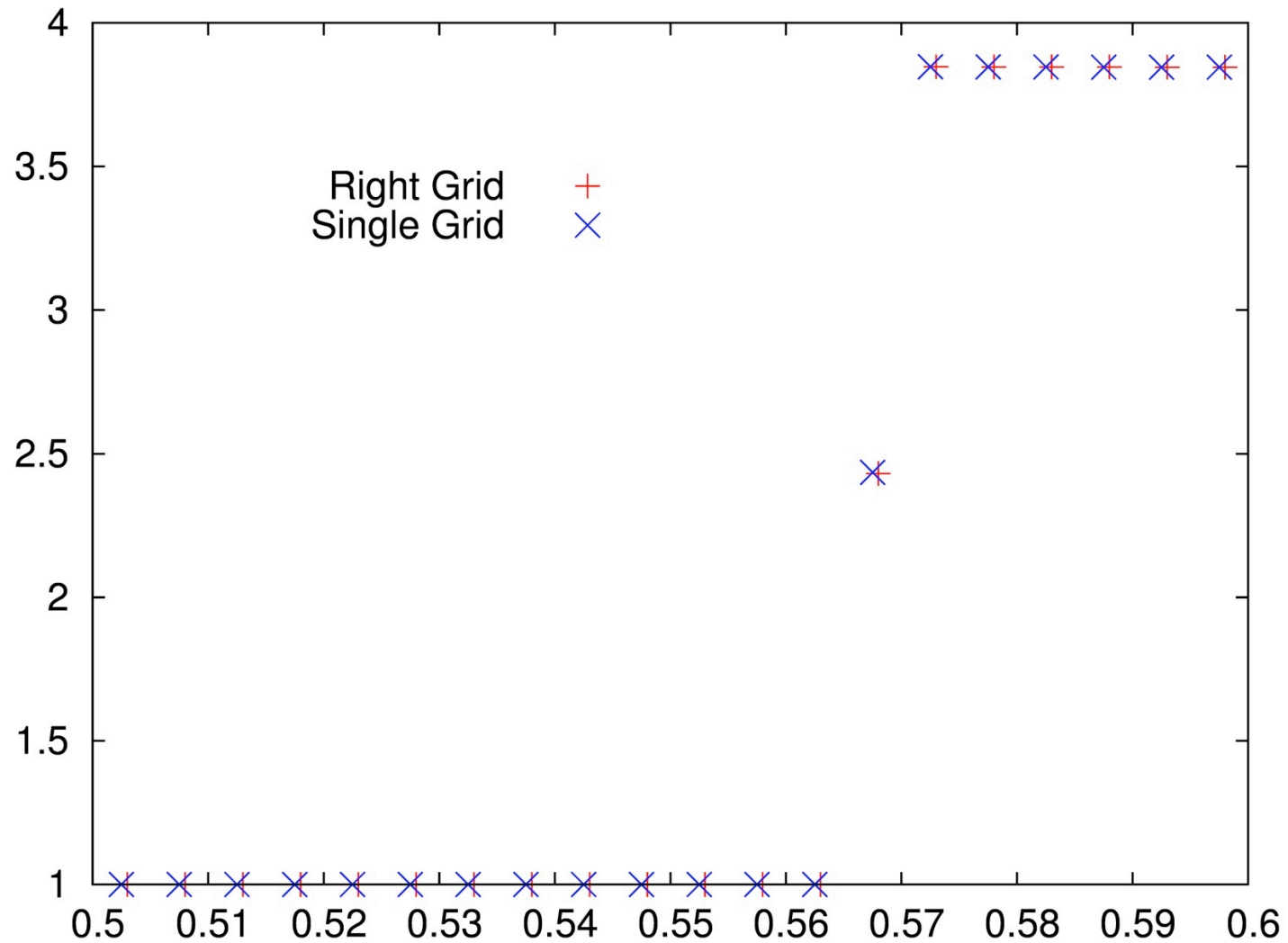
Slow Moving Shock (Euler)



Single Grid Comparison

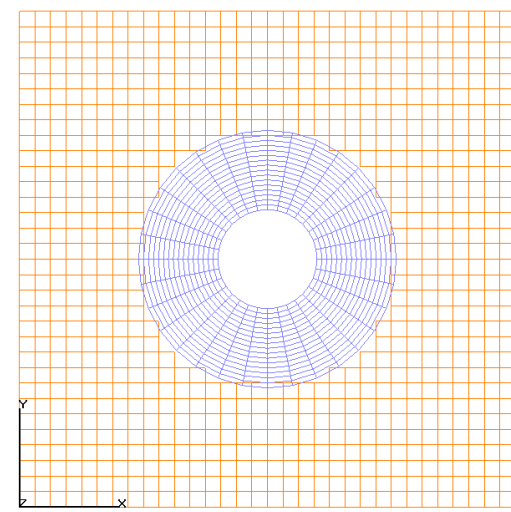
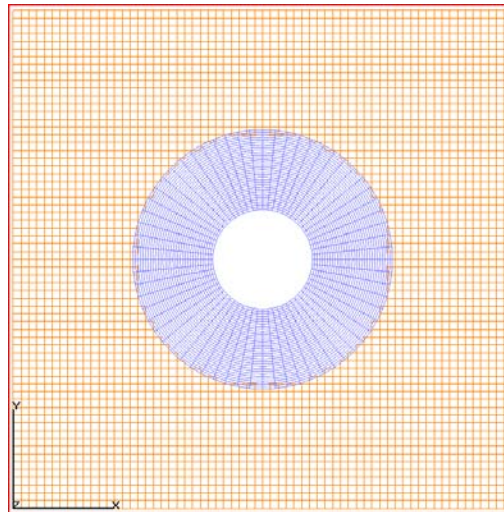
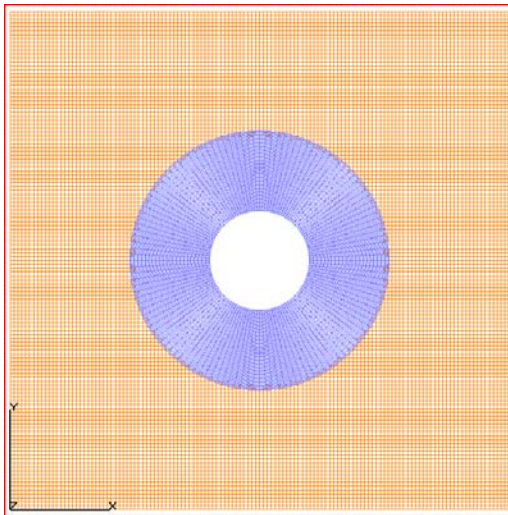


Single Grid Comparison



Conclusions

- Within the context of this framework strict lack of conservation becomes merely another source of numerical error to be removed by grid refinement
- Convergence to a weak solution is **guaranteed** as grid independence is achieved
- Method should be of some benefit to
 - Multigrid calculations
 - Grid convergence studies
 - ❖ No need to maintain large overlap in finest grids



Conclusions

- Need to develop algorithms for determining minimal overlap
- Achieving higher-order (> 3) accuracy problematic
 - Although a DG based chimera method appears to fit nicely within this framework
 - The approach can be generalized to finite difference based codes by using a “finite-point” method at the interfaces

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That's All Folks

- Questions
- Constructive Criticism